
Flavour physics – Lecture 1

Jim Libby (IITM)

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Outline

- What is flavour physics?
- Some theory and history
- CKM matrix

What is flavour physics?

Fermions ("matter")	Bosons ("forces")
$ \left\{ \begin{array}{l} \text{Quarks} \\ uu\bar{u} \quad cc\bar{c} \quad tt\bar{t} \\ dd\bar{d} \quad ss\bar{s} \quad bb\bar{b} \\ \\ \text{Leptons} \\ e \quad \mu \quad \tau \\ \nu_e \quad \nu_\mu \quad \nu_\tau \end{array} \right\} \times \left\{ \begin{array}{l} \text{MATTER} \\ \text{ANTIMATTER} \end{array} \right\} $	$ \begin{array}{l} gggggggg \\ \gamma \\ W^+ \\ W^- \\ Z \\ \\ H \end{array} $

SM parameters

- 3 gauge couplings
- 2 Higgs parameters
- 6 quark masses
- 3 quark mixing angles and 1 phase
- 3 + [3] lepton masses
- [3 lepton mixing angles and 1 phase]

[with massive neutrinos]

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[with massive neutrinos]

FLAVOUR PARAMETERS

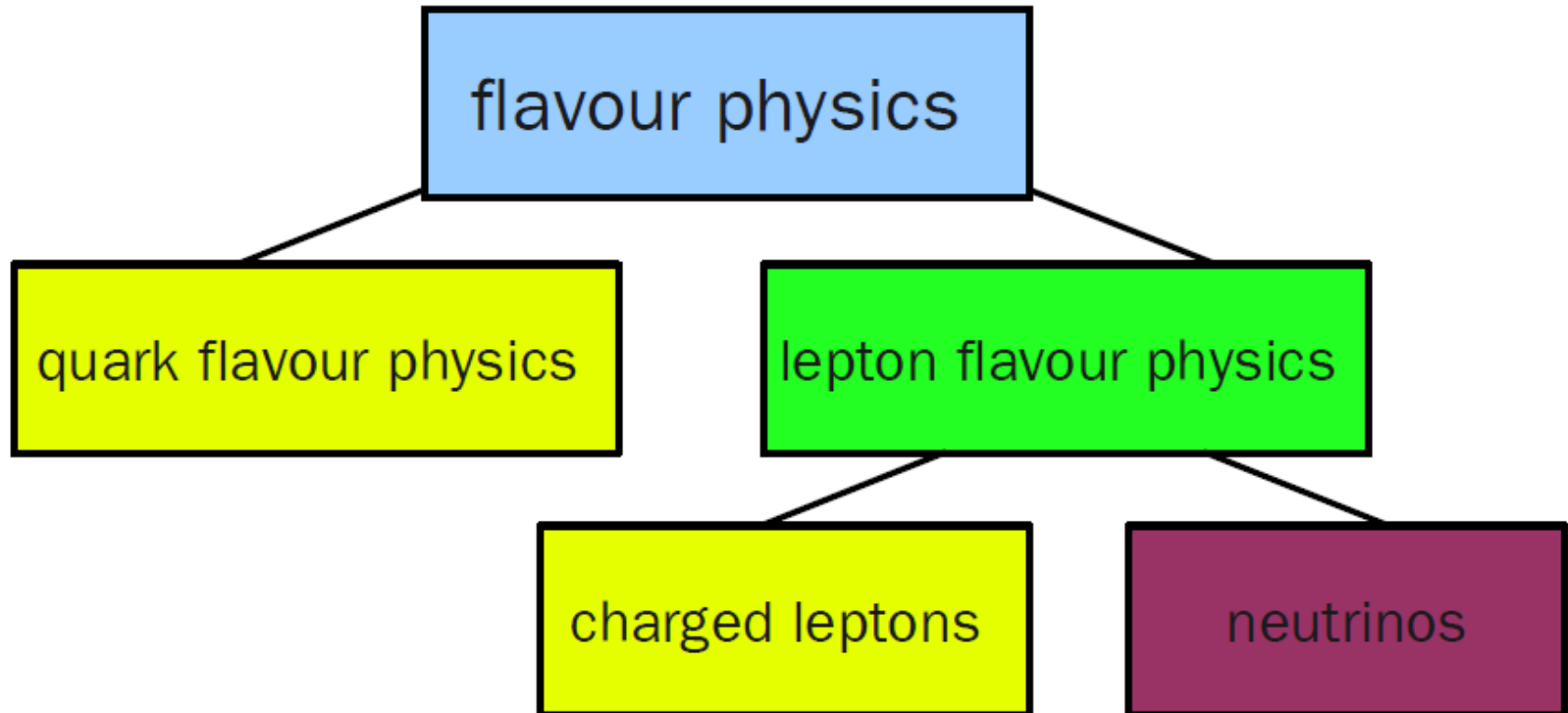
What is not flavour physics?

- QCD
 - But important for interpretation of flavour variables
- Electroweak physics
 - But linked to weak couplings and absence of flavour changing neutral currents
 - Higgs mechanism generates quark masses and mixings
- High- p_t – searches for new heavy particles
 - Complementary effects in virtual processes in flavour physics

Puzzles in flavour physics

- Why three generations of quarks and leptons
- Mass hierarchies of quarks and leptons
- Hierarchy of weak mixing of quarks
 - What about leptons?
- Charged lepton coupling universality
- Absence of flavour changing neutral currents
- Symmetry principles and their violation
 - P and C
 - CP and T
 - Baryon and anti-baryon asymmetry in the universe
 - Lepton flavour violation?

Flavour organigram



A few words about neutrinos...

- Parity violation

- Neutrinos are left handed (chiral) states

- Antineutrinos are righthanded

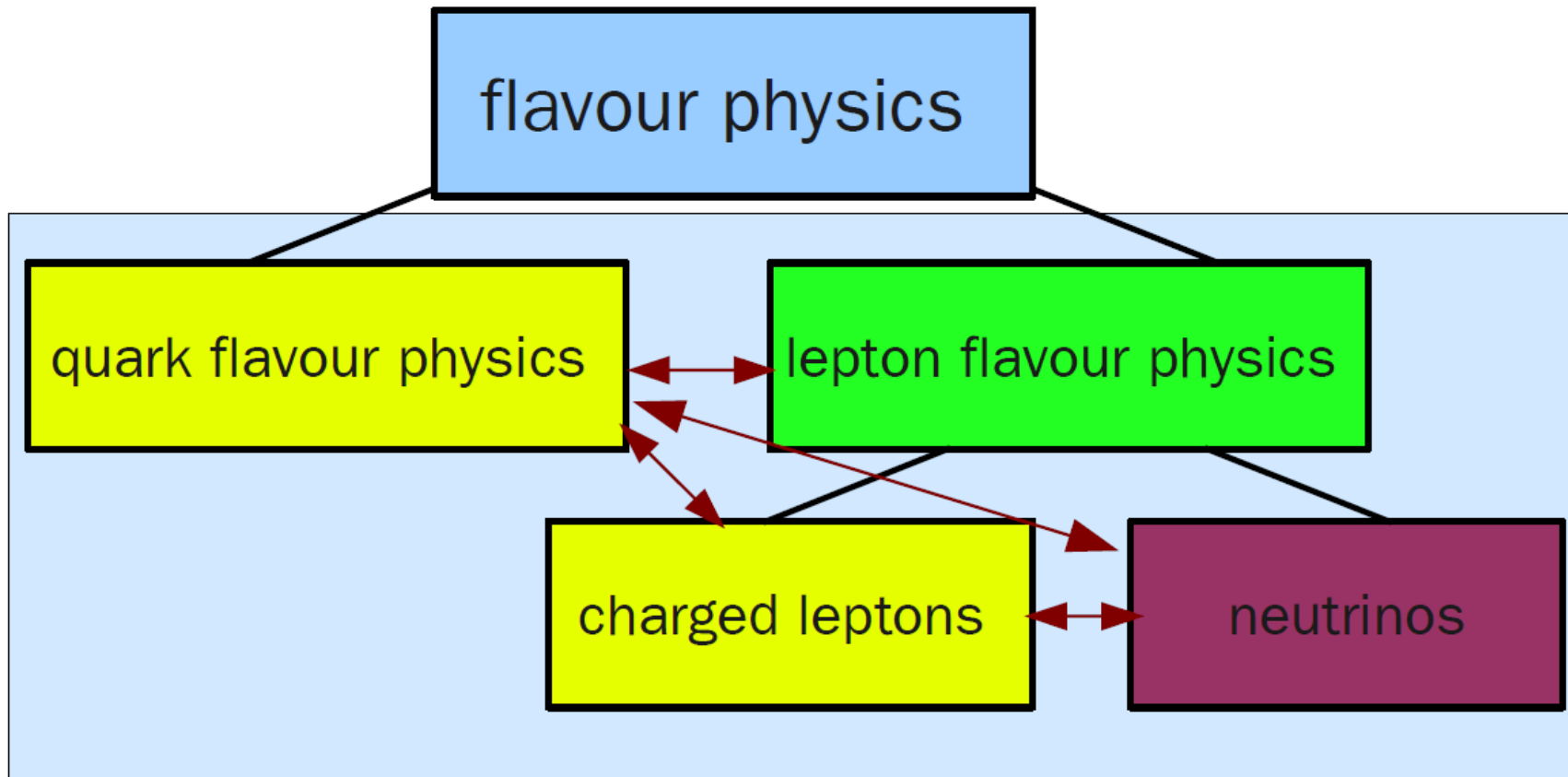
- But not massless

- So where are right handed neutrinos?

- Massive field with many intriguing questions

-but this topic will not be covered in depth here as we will deal with more traditional discussion of flavour in the hadronic weak sector

Ultimate goal – unified description of interplay between the different divisions



Alternatively can think about a division

1. Flavour changing physics
 - Decays and couplings
 - Lifetimes
 - Mixing
 - CP violation
2. Flavour conserving physics
 - Masses
 - Dipole moments

A historical review

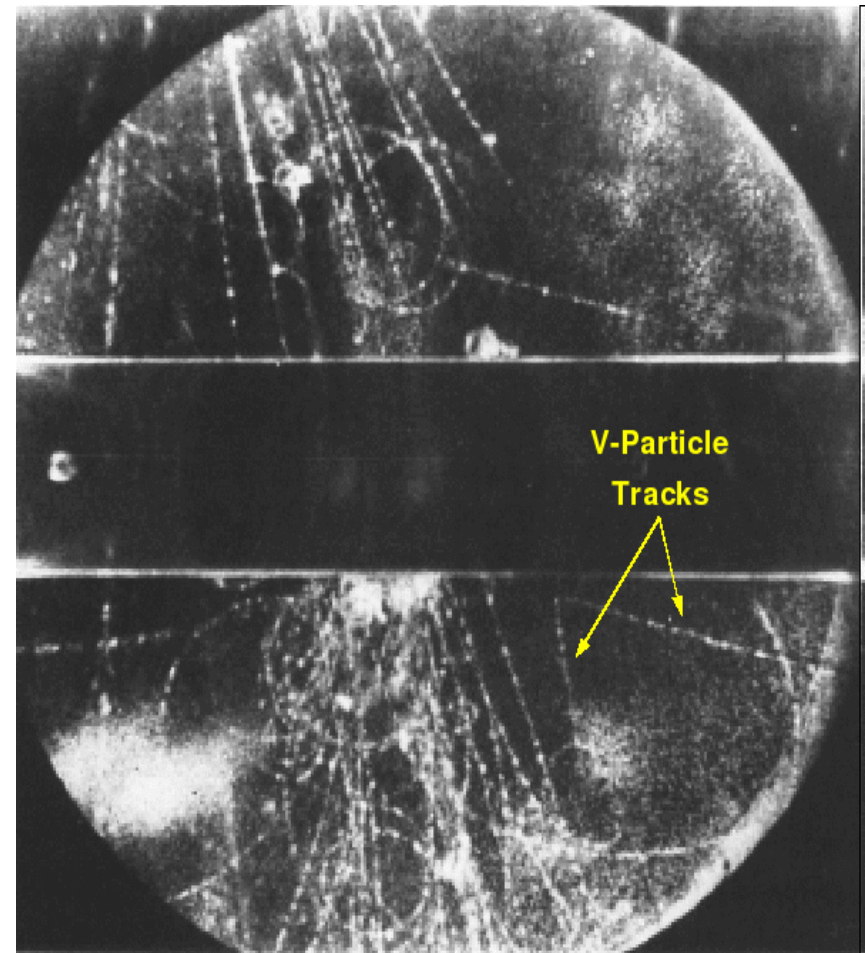
- Isospin
- Strangeness
- Quark model
- τ - θ puzzle
- P and C violation
- Mixing and GIM mechanism
- CP violation
- Sakaharov and the link to cosmology

Isospin

- Proton and neutron different charge but
 - Near identical mass
 - Strong couplings
- Heisenberg proposed new quantum number isospin (I) of which n and p form a doublet $I=1/2$,
 - $I_3(p)=1/2$, and $I_3(n)=-1/2$
- Later observed that pions form a triplet $I=1$
 - $I_3(\pi^-)=-1$, $I_3(\pi^0)=0$, and $I_3(\pi^+)=+1$
- Link between symmetry and invariance
 - n and p strong interactions invariant under $SU(2)$ rotation
 - Works because u and d almost degenerate in mass
- Isospin is a way of saying the strong interaction is flavour blind and the third component counts upness and downness flavour

Strangeness

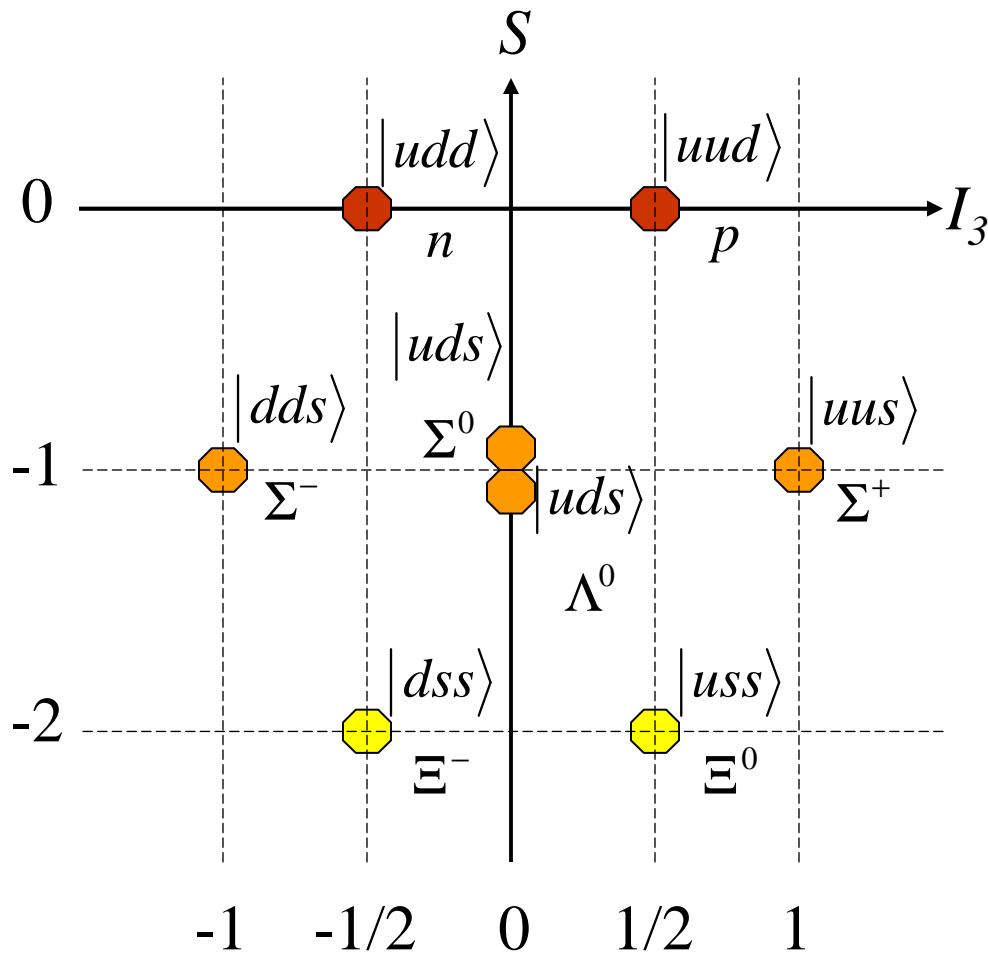
- Rochester and Butler (1947)
 - Neutral particle in cosmic rays decay to two charged pions
 - forms a V
 - Long-lifetime - strange
 - Also, saw evidence of charged particle that decayed to a pion and something else
- Pais (1952) introduced new quantum number strangeness
 - Conserved in strong interactions
 - Violated in weak interactions
 - **Flavour change**



Quark model – Gell-Mann (1961-1964)

- All very comfortable with quarks now but only accepted after discovery of charm in 1974
 - More later
- Neatly explained the zoo of particles discovered in 1950s and 1960s
- The $SU(3)$ flavour symmetry – eightfold way - gave rise to groupings
 - Baryon octet and decuplet
 - Meson nonet

Baryon octet



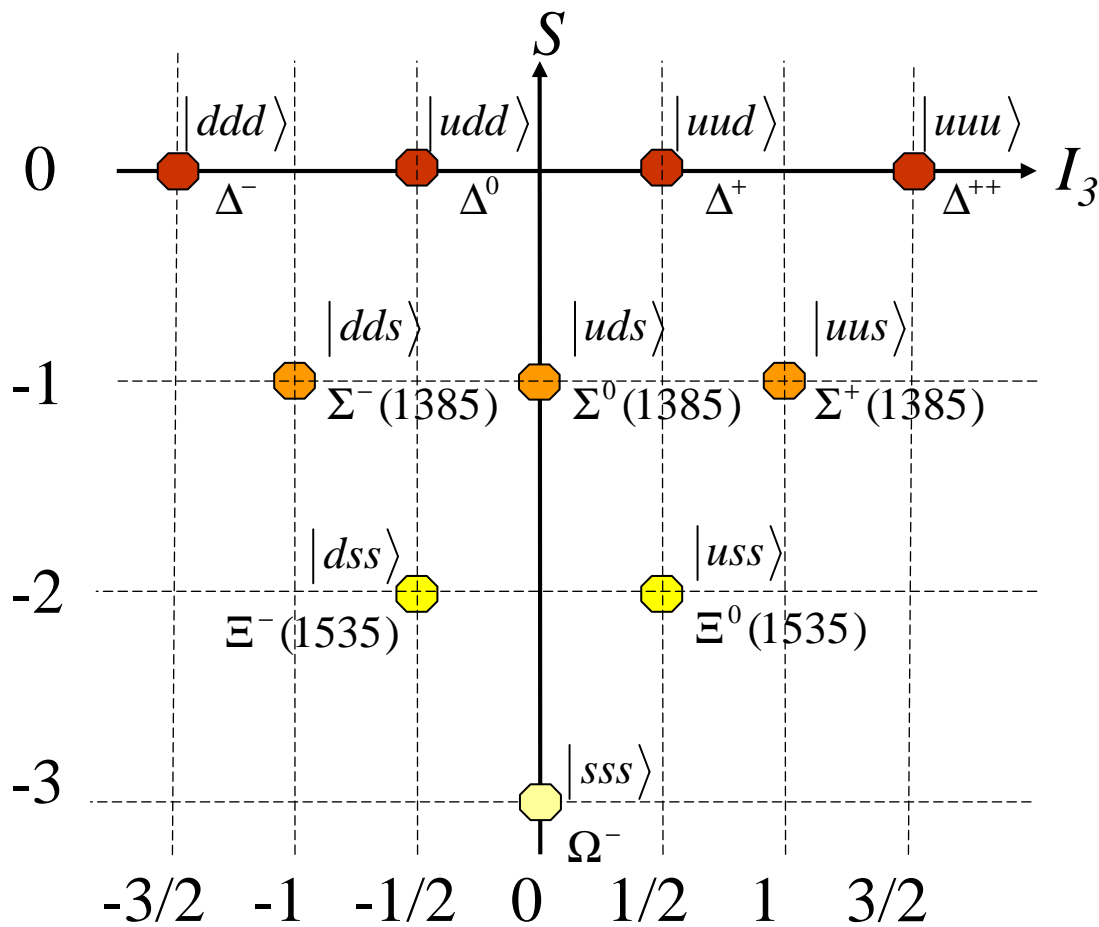
Spin-parity $J^P = 1/2^+$

- Two quark spins aligned the third antiparallel
- Parity from +ve intrinsic parity of quarks

S and I_3 are additive quantum numbers

- S = strangeness ($S = -1$ for a strange quark and $+1$ for an anti-strange quark)
- I_3 = third component of isospin ($I_3 = +1/2$ for up quarks and $-1/2$ for down quarks)

Baryon decuplet



Spin-parity $J^P=3/2^+$

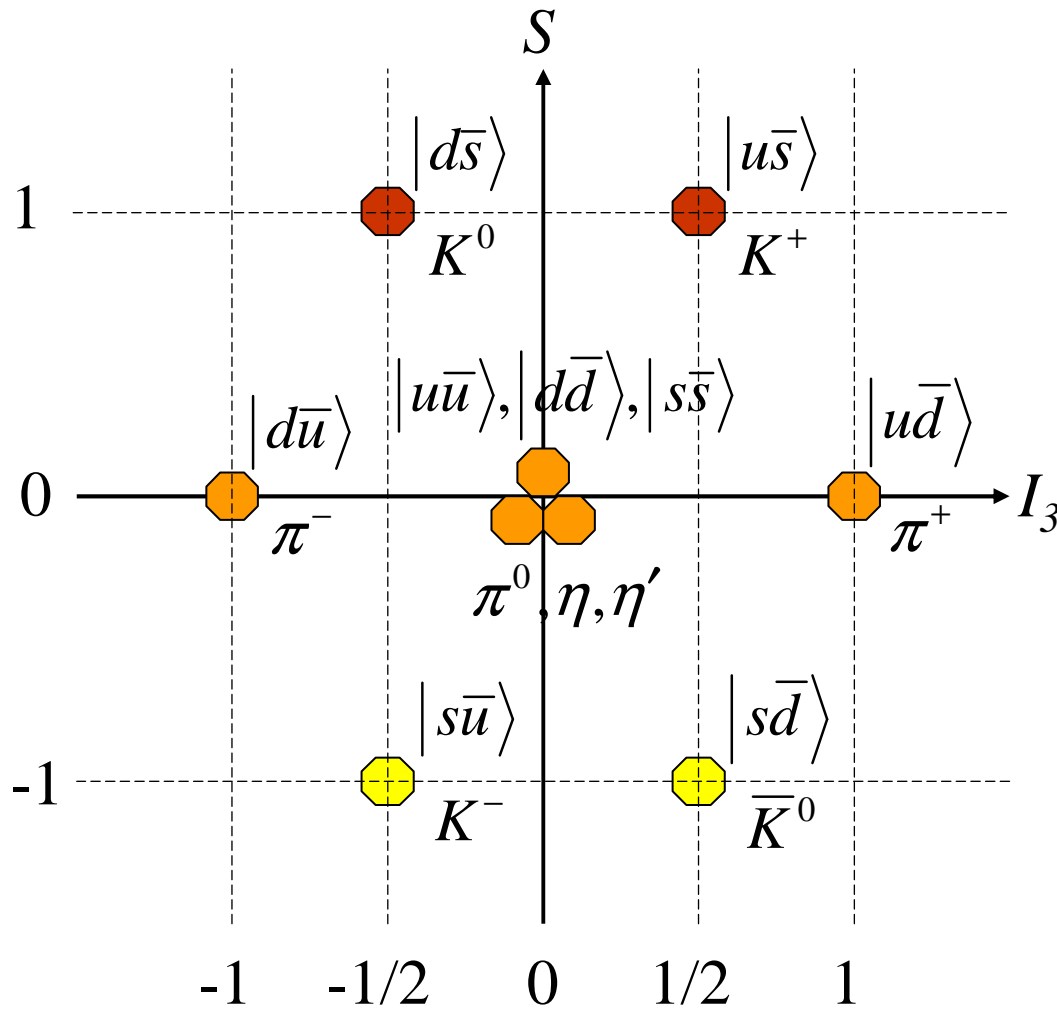
All quark spins aligned

Quark compositions

$|qq'q''\rangle$ are not the same as quark wavefunctions

The discovery of 1964 of Ω^- predicted by the model first vindication

Meson nonet: $J^P=0^-$



Spin-parity $J^P=0^-$

Quark spins antiparallel

Pseudoscalars: zero spin (scalar quantity) but antisymmetric under parity operation

The three $S=I_3=0$ states quark wavefunctions are different admixtures of

$|u\bar{u}\rangle, |d\bar{d}\rangle$ and $|s\bar{s}\rangle$

The θ - τ problem

- Spin 0 θ particle decayed to $\pi^+\pi^0$
 - Even parity final state
- Spin 0 τ particle decayed to $\pi^+\pi^-\pi^0$
 - Odd parity final state
- But masses the and lifetimes the same
- Parity violation discovered in 1957
 - Wu (after Lee and Yang's suggestion)
- θ and τ the same – K^0
 - Parity violated maximally in flavour changing weak decays

P, C and CP

- P maximal violation in weak processes
 - All neutrinos lefthanded
- But C also violated in weak decays
 - No righthanded neutrinos
- But product CP assumed to be a conserved quantity (Landau 1957)
 - So CP distinguishes matter from antimatter
- Note that CP and T combined is conserved in all Lorentz invariant gauge field theories (i.e. SM)
 - Will return to this when CP violation is discussed

Weak couplings - Cabibbo (1963)

- When corrected for phase space it was found that $s \rightarrow u$ weak transition rates suppressed by a factor of ~ 20 compared to $d \rightarrow u$
 - $K^+ \rightarrow \mu^+ \nu$ compared to $\pi^+ \rightarrow \mu^+ \nu$
- Also, small differences in G_F measured in $d \rightarrow u$ transitions compared to muon decay
- Above could be explained if flavour eigenstates and weak eigenstates different

$$\begin{pmatrix} d \\ s \end{pmatrix}_{weak} = \begin{bmatrix} \cos \theta_C & \sin \theta_C \\ -\sin \theta_C & \cos \theta_C \end{bmatrix} \begin{pmatrix} d \\ s \end{pmatrix}_{flavour} \quad \sin \theta_C \approx 0.22 \quad \cos \theta_C \approx 0.98$$

Neutral kaon mixing

- Consequence of strangeness and weak states not being the same is that the physical states of neutral kaons are almost equal admixtures of the strangeness eigenstates



$$K_S^0 \approx \frac{1}{\sqrt{2}}(K^0 + \bar{K}^0) \quad K_L^0 \approx \frac{1}{\sqrt{2}}(K^0 - \bar{K}^0)$$

- K_S is CP-even (decays to two pions)
- K_L is CP-odd (decays to three pions and semileptonically)

GIM mechanism

- There was a puzzle
 - $K^+ \rightarrow \mu^+ \nu$ so why not $K^0 \rightarrow \mu^+ \mu^-$
 - $K^+ \rightarrow \pi^0 \mu^+ \nu$ so why not $K^0 \rightarrow \pi^0 \mu^+ \mu^-$
- GIM (Glashow, Iliopoulos and Maiani) mechanism (1970)
 - Leads to no flavour changing neutral currents (FCNC)
 - Suppression via loops
 - Requires that quarks come in doublets
 - Predicted a fourth quark - charm

GIM continued

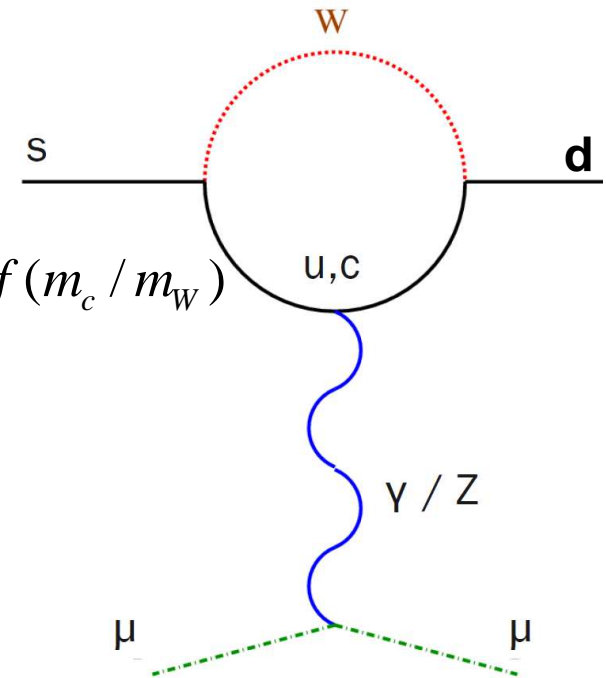
$$\begin{pmatrix} u & c \end{pmatrix}_{flavour} \begin{bmatrix} \cos \theta_C & \sin \theta_C \\ -\sin \theta_C & \cos \theta_C \end{bmatrix} \begin{pmatrix} d \\ s \end{pmatrix}_{flavour}$$

$$A(s \rightarrow d \mu^+ \mu^-) = \cos \theta_C \sin \theta_C f(m_u / m_W) - \cos \theta_C \sin \theta_C f(m_c / m_W)$$

As $m_u \ll m_W$ and $m_c \ll m_W$

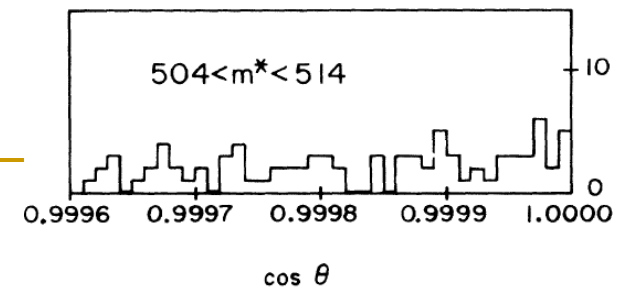
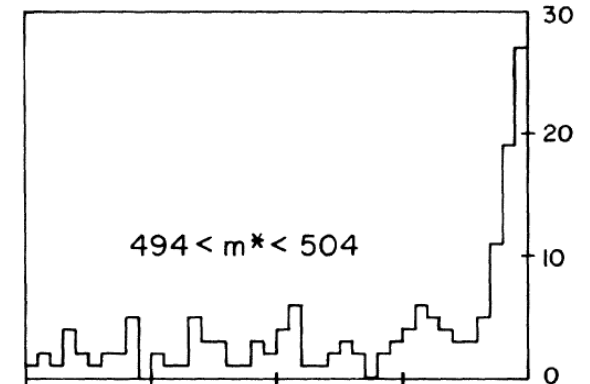
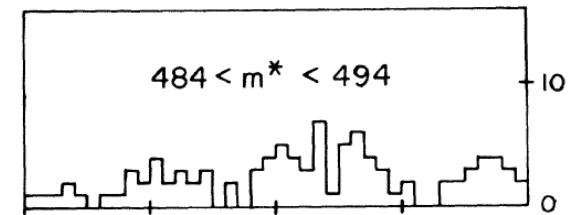
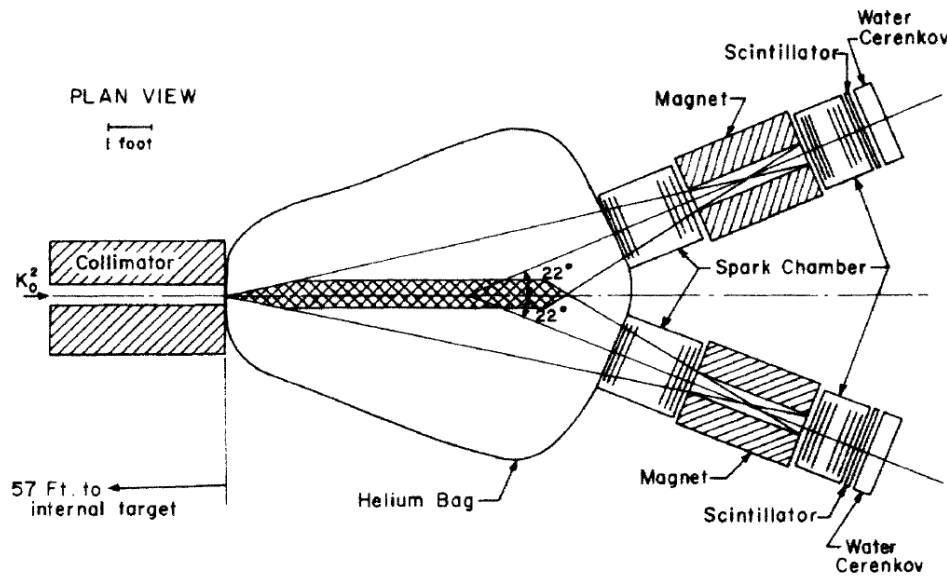
$A \sim 0$

The small rate at which these processes occur allowed value of charm mass to be predicted



CP violation

- Another puzzle was the observation in 1964 of $K_L \rightarrow \pi\pi$ by Cronin, Fitch et al.
- This violates CP!



Sakharov conditions – CP and cosmology

- Proposed by A. Sakharov (1967)
- N_B observed a billion times bigger than expected
- No antimatter observed
- Three conditions for evolution of matter dominated universe from symmetric initial state
 1. Baryon number violation
 2. C and CP violation
 3. Lack of thermal equilibrium

CKM matrix

**Predicted 3rd generation even
before charm had been observed**

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Progress of Theoretical Physics, Vol. 49, No. 2, February 1973

***CP*-Violation in the Renormalizable Theory of Weak Interaction**

Makoto KOBAYASHI and Toshihide MASKAWA

Department of Physics, Kyoto University, Kyoto

(Received September 1, 1972)

In a framework of the renormalizable theory of weak interaction, problems of *CP*-violation are studied. It is concluded that no realistic models of *CP*-violation exist in the quartet scheme without introducing any other new fields. Some possible models of *CP*-violation are also discussed.

CKM matrix

- Extension of Cabibbo's two by two mixing matrix
- Unitary complex matrix
 - 18 parameters
- Unitary constraints ($VV^\dagger=I$)
 - 9 parameters
- Quark fields can absorb five unobservable phases
 - 4 parameters
 - 3 mixing angle and 1 phase

$$(u \quad c \quad t) \begin{bmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{bmatrix} \begin{pmatrix} d \\ s \\ b \end{pmatrix}$$

**Responsible for
CP violation**

Summary of lecture 1

- Flavour physics integral to development of SM
 - Mixing
 - FCNC
 - CP violation and CKM
- Led to prediction, and as we shall see, properties of the charm, bottom and top quark
 - Low energy probe of high energy scales