

# Rapportuer -II (Theory)

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## ICPAQGP Series:

- 7<sup>th</sup> ICPAQGP - 2015 (Kolkata)
- 6<sup>th</sup> ICPAQGP - 2010 (Goa)
- 5<sup>th</sup> ICPAQGP - 2006 (Kolkata)
- 4<sup>th</sup> ICPAQGP - 2001 (Jaipur)
- 3<sup>rd</sup> ICPAQGP - 1997 (Jaipur)
- 2<sup>nd</sup> ICPAQGP - 1993 (Calcutta-Then)
- 1<sup>st</sup> ICPAQGP - 1988 (Bombay - Then)

# Beginning

- In the beginning, the search for QGP in high energy nuclear collision was very much like an exploratory endeavor
  - ☞ As in most explorations, the enthusiasm of the explorers was greater than their knowledge of what to expect or look for!!
  - ☞ General feeling was that if QGP would be produced by any mechanism, it would manifest in a variety of unknown but dramatic ways!!!
  - ☞ Something like an order parameter, which doesn't change over a period but changes suddenly at a point!

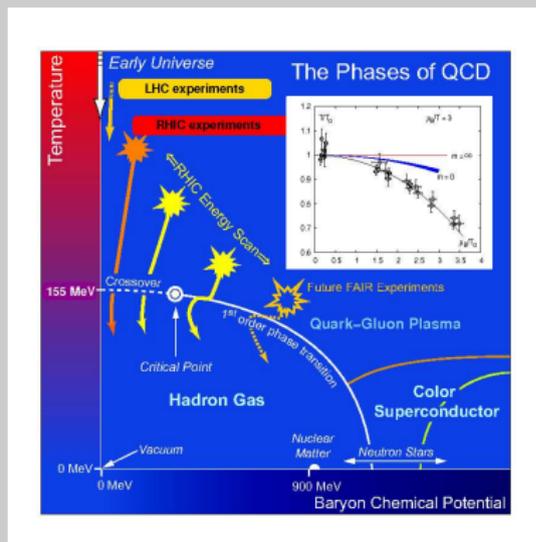
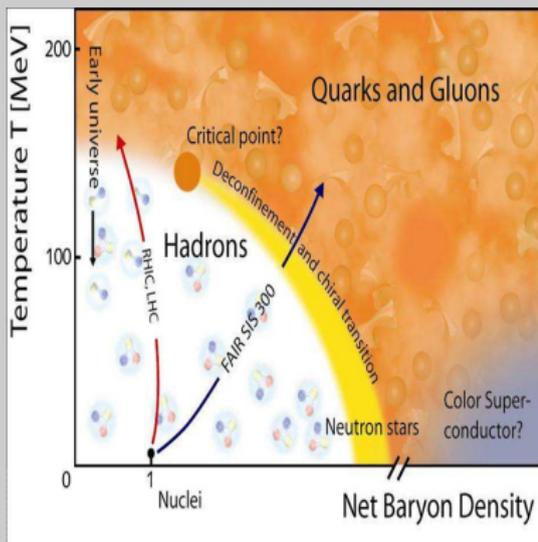
## Aim of Heavy-Ion Collisions Experiments:

- 🎯 Accelerate heavy stable ions with energy as much as one can and then collide them travelling at relativistic speed
- 🎯 Not to focus on energy but on energy density created
- 🎯 Not to focus on fine/precision physics but on **collective** physics
  - 👉 Explores properties of matter under extreme conditions
  - 👉 Much focus on **Quark-Gluon Plasma** formation: the primordial form of matter that existed in the universe shortly after the Big Bang
  - 👉 Study how **QCD** works in unusual conditions

# Heavy-ion Collisions Experiments

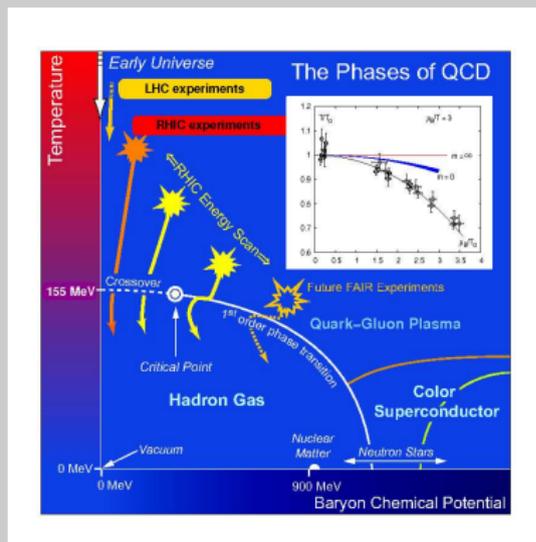
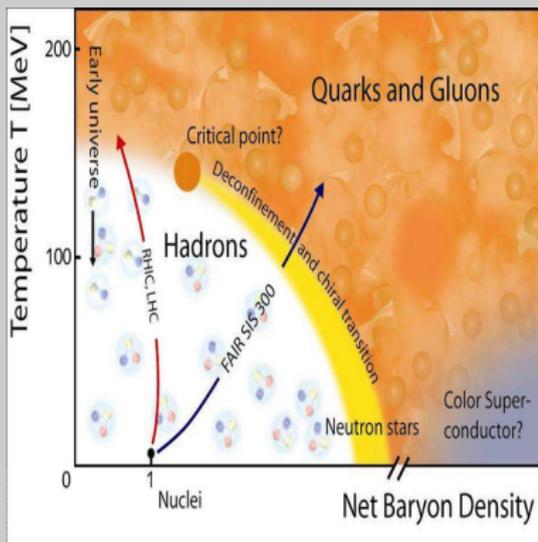
- AGS (1986-1996):** Si, Au;  $\sqrt{s} = 2.5 - 5.5$  GeV/A;  
 Barely reached the energy density required for QGP formation
- SPS:** O, Si, Pb; 158 GeV, 200 GeV  
 Indication of a possible onset of deconfinement!
- RHIC (2000-...):** d, 3He, Cu, Au, U, d;  $\sqrt{s} = 7.7 - 200$  GeV/A  
 Evidence of Deconfinement or partonic degrees
- LHC(2010-...):** Pb-Pb, p-Pb;  $\sqrt{s} = 2.76, 5.5$  ATeV  
 Evidence of Deconfinement or partonic degrees
- FAIR:** Expected in 2018 !! ;  $E_{\text{lab}} = 10(35)$  GeV/A  
 ( $\sqrt{s} < 4.5(8.4)$  GeV for Au/U)

# Present-day understanding of the QCD phase diagram



- Theory (QCD and models based on it)
- Numerical Experiments (LQCD)
- Astrophysical observations
- Mini Bang (Laboratory Expt.)

# Present-day understanding of the QCD phase diagram



- Theory (QCD and models based on it)
- Numerical Experiments (LQCD)
- Astrophysical observations
- Mini Bang (Laboratory Expt.)
- Complex: Difficult to span the whole PD by a single effort !

# Present-day understanding of the QCD phase and/or diagram

- Higher the collision energy the closer is the created system located near temperature axis
- A crossover at RHIC and LHC; No true phase transition; LQCD also exhibited.
- Preliminary to attribute a definite position for critical point in phase diagram (BES-I at RHIC and LQCD)
- Data appear to demand an explanation beyond a purely hadronic scenario  $\Rightarrow$  non-hadronic source (partonic degrees)
- Collective behaviours  $\Rightarrow$  Circumstantial evidences

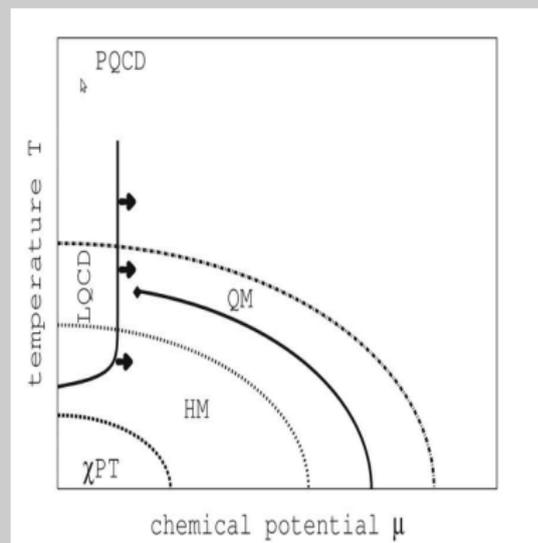
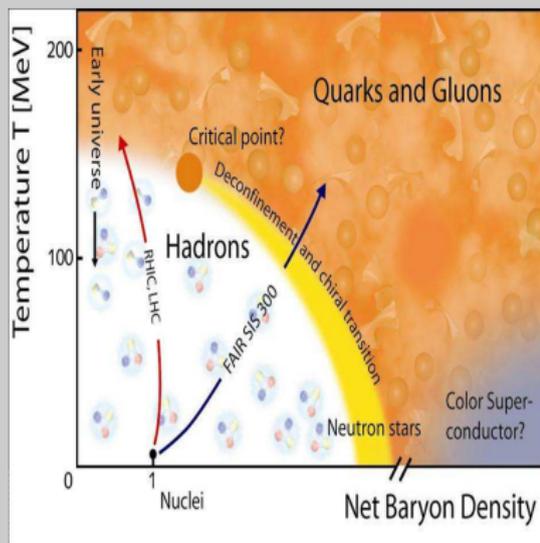
# Evidences of Partonic Degrees of Freedom at RHIC and LHC

- Hot and dense matter produced is a liquid (with low  $\eta/s$ ) in which partons flow individually; strongly coupled near perfect fluid
- QGP produced in RHIC is (on average) closer to perfection than that at LHC.

Large opacity caused energy-loss that led to suppression of  $R_{AA}$  (for both Light Heavy hadrons)

- Light quarks (u, d, s) completely thermalize and recombine to hadrons
- Melting of Quarkonia States and regeneration.
- New data is a joy to experimentalists but a hard task for Theoreticians to explain them!

# Theoretical Approaches

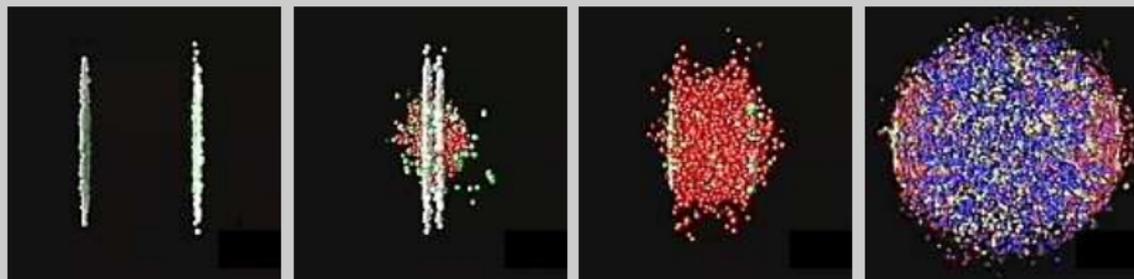


- All regions of the PD by the first principle QCD calculations !  
**Not Yet!!**
- Interface of Nuclear Physics and Particle Physics

# Interface of particle physics & high-energy nuclear physics

- Draws heavily from QCD: pert, non-pert, semiclassical
- Overlaps with:
  - ☞ Thermal Field Theory
  - ☞ Relativistic Fluid Dynamics
  - ☞ Kinetic or Transport theory
  - ☞ Quantum Collision Theory
  - ☞ String Theory
  - ☞ Statistical Mechanics & Thermodynamics

# Quark-Gluon Plasma: A new phase of QCD



QGP ➤ Expands ➤ Hadronizes (Detector)

Expansion: Hydrodynamics (ideal, viscous) + Initial Conditions

Two Aspects

👉 **Snap Shot properties** (Given Temp. & Chem. Pot<sup>l</sup>)

👉 **Experimental Information**

# Quark-Gluon Plasma: A new phase of QCD

## 📌 Snap Shot:

👉 Thermodynamic Quantities } **Hydro inputs**  
 ( $F, P, \dots, \text{EoS}$ )

👉 Transport Coefficients } **Hydro/FP Eq. inputs**  
 ( $\sigma, \gamma_d, \mathcal{D}, \eta/s$ )

👉 Various Susceptibilities  
 ( $\chi_c, \chi_q, \dots, \text{response}$ )

👉 Screening of Plasma

👉 Interactions

(Coll., Rad., E-loss)

👉 Particle Production Rates

( $l^+l^-, \gamma, \dots$ )

## 📌 Exptl. Data:

➡ Fluctuations, Critical Point

➡ Binary states:  $J/\psi, \Upsilon$

➡ Hadron spectra;

$R_{AA}, v_n$

➡  $l^+l^-, \gamma$  spectra

# ICPAQGP-2015

- 👉 Stdents' Day (1/2/2014)
- 👉 Overview talks: 4
- 👉 Plenary Talks: 30
- 👉 Parallel sessions: 65
- 👉 Posters: 40

⇒ ICPAQGP2015 has provided extensively the field's history and developments until date

# Rapportuer-I

## Rapportuer-I by Raju Venugopalan

- Initial States (pp, pA, AA)
- Hydrodynamics
- Thermalization
- Final states of (pp & pA)
- Flow coefficients & extraction of  $\eta/s$
- Quarkonic Phase
- Photons (!)

# Rapportuer-II

## Rapportuer-II by M G. Mustafa

- ☞ Thermodynamics, Susceptibilities & CP
- ☞ Transport Coefficients
- ☞ Energy Loss and Jets
- ☞ Dileptons
- ☞ Quarkonia Melting
- ☞ Astrophysical aspects

# Students' Day

- Stage for this conference was set by a series of pedagogical lectures
  - Signals of Quark-Gluon Plasma by Jan-e Alam
  - QGP in the Universe - from the early universe to the present by Ajit M Srivastava
  - Relativistic Fluid Dynamics and Collective Flow in HIC by Rajeev Bhalerao
  - Thermal Model by Jean Cleymans
  - Heavy-ion Physics (past/present/future) by Jurgen Schukraft
  - Compressed Baryonic Matter by Subhasis Chattopadhyay
  - Physics with Electron-Ion Collider by Abhay Deshpande

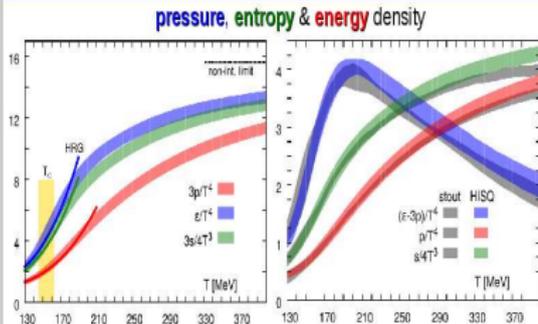
# Thermodynamics, Cumulant and Critical Point

-  Precise knowledge of equation of state (pressure) of QCD at high density and temperature has important significance for the analysis of HIC experiments.

# Thermodynamics, Cumulant and Critical Point

- LQCD EOS at  $\mu = 0$  [F. Karsch, P. Hegde]

## Equation of state of (2+1)-flavor QCD

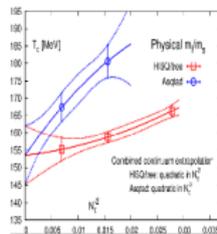


A. Bazavov et al. (hotQCD),  
Phys. Rev. D90 (2014) 094503

- improves over earlier hotQCD calculations:  
A. Bazavov et al., Phys. Rev. D80, 014504 (2009)
- consistent with results from Budapest-Wuppertal (stout): S. Borsanyi et al., PL B730, 99 (2014)

– up to the crossover region the QCD EoS agrees quite well with hadron resonance gas (HRG) model calculations; **However**, QCD results are systematically above HRG

## Equation of state and transition temperature



$$T_c = (154 \pm 9) \text{ MeV}$$

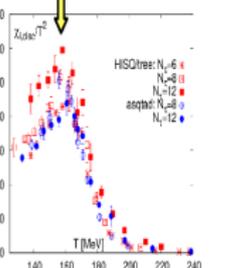
- well defined pseudo-critical temperature
- quark mass dependence of susceptibilities consistent with  $O(4)$  scaling

A. Bazavov et al. (hotQCD),  
Phys. Rev. D85, 054503 (2012), arXiv:1111.1710

lattice:  $N_\sigma^3 \cdot N_\tau$   
temperature:  $T = 1/N_\tau a$

Critical temperature from location of peak in the fluctuation of the chiral condensate (order parameter):

$$\chi_I = \frac{T \partial^2 \ln Z}{V \partial m_l^2} = \chi_{I,disc} + \chi_{I,con}$$



consistent with Y. Aoki et al. JHEP 0906 (2009) 088

# Thermodynamics, Cumulant and Critical Point

- LQCD at  $\mu \neq 0$  [F. Karsch, P. Hegde]

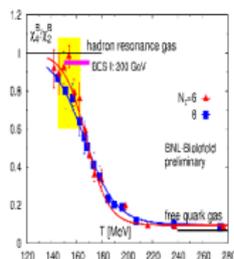
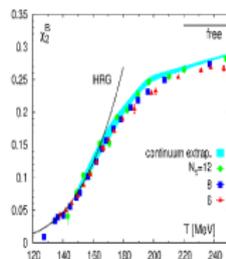
Equation of state of (2+1)-flavor QCD:  $\mu_B/T > 0$

$$\frac{\Delta(T, \mu_B)}{T^4} = \frac{P(T, \mu_B) - P(T, 0)}{T^4} = \frac{\chi_2^B}{2} \left( \frac{\mu_B}{T} \right)^2 \left( 1 + \frac{1}{12} \frac{\chi_4^B}{\chi_2^B} \left( \frac{\mu_B}{T} \right)^2 \right)$$

variance of net-baryon  
number distribution

kurtosis\*variance

$\kappa_B \sigma_B^2$



F. Karsch, ICQAQGP 2015

# Thermodynamics, Cumulant and Critical Point

- LQCD at  $\mu \neq 0$  [F. Karsch, P. Hegde]

## Taylor expansion of the pressure and critical point

$$\frac{P}{T^4} = \sum_{n=0}^{\infty} \frac{1}{n!} \chi_n^B(T) \left(\frac{\mu_B}{T}\right)^n$$

for simplicity:  $\mu_Q = \mu_S = 0$

estimator for the radius of convergence:

$$\left(\frac{\mu_B}{T}\right)_{\text{crit},n}^X \equiv r_n^X = \sqrt{\frac{n(n-1)\chi_n^B}{\chi_{n+2}^B}}$$

- radius of convergence corresponds to a critical point only, iff

$$\chi_n > 0 \text{ for all } n \geq n_0$$

forces  $P/T^4$  and  $\chi_n^B(T, \mu_B)$  to be monotonically growing with  $\mu_B/T$

$$\text{at } T_{CP}: \kappa_B \sigma_B^2 = \frac{\chi_4^B(T, \mu_B)}{\chi_2^B(T, \mu_B)} > 1$$

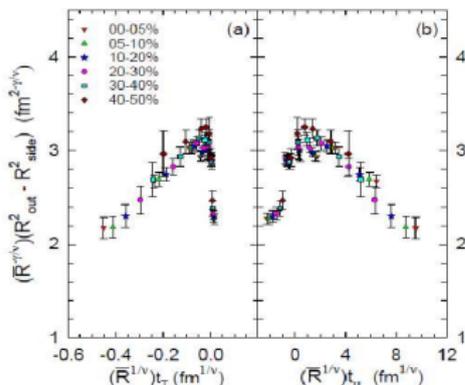
if not:

- radius of convergence does not determine the critical point

- Taylor expansion can not be used close to the critical point

$$R^{-\gamma/\nu} \times (R_{\text{out}}^2 - R_{\text{side}}^2) \text{ vs. } R^{1/\nu} \times t_T, \quad t_T = (T - T^{\text{exp}})/T^{\text{exp}}$$

$$R^{-\gamma/\nu} \times (R_{\text{out}}^2 - R_{\text{side}}^2) \text{ vs. } R^{1/\nu} \times t_{\mu_B}, \quad t_{\mu_B} = (\mu_B - \mu_B^{\text{exp}})/\mu_B^{\text{exp}}$$



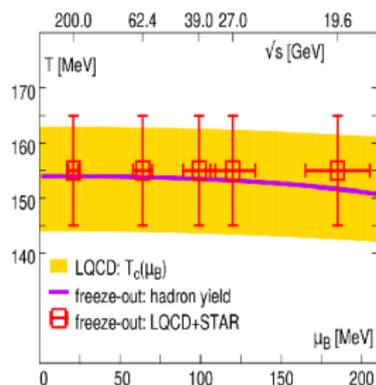
**\*\*\*A Clear initial indication for the CEP\*\*\***

Roy A. Lacey, Stony Brook University, ICQAQGP, Kolkata, India

# Thermodynamics, Cumulant and Freeze out

- LQCD [S. Mukherjee]; Model Calculation [K. Redlich]

## Freeze-out at RHIC

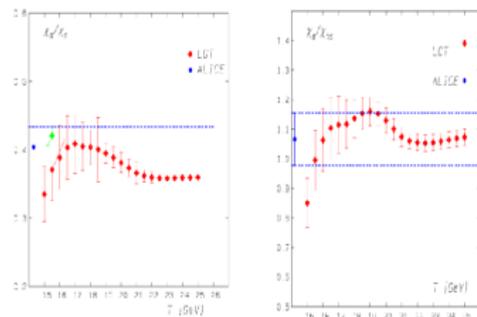


freeze-out in RHIC takes place close to the phase boundary

freeze-out from hadron yield parametrization:  
Andronic et.al., J. Phys. G38, 124081 (2011),  
with  $T(\sqrt{s} \rightarrow \infty) = 154 \text{ MeV}$

need better  $T^{\dagger}$  determination

## Consider T-dependent LQCD ratios and compare with ALICE data



- The LQCD susceptibilities ratios are weakly T-dependent for  $T \geq T_c$
- We can reject  $T \leq 0.15 \text{ GeV}$  for saturation of  $\chi_4, \chi_6$  and  $\chi_{Q5}$  at LHC and fixed to be in the range  $0.15 < T \leq 0.21 \text{ GeV}$ , however
- LQCD  $\Rightarrow$  for  $T > 0.163 \text{ GeV}$  thermodynamics cannot be anymore described by the hadronic degrees of freedom

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# Take Home from LQCD

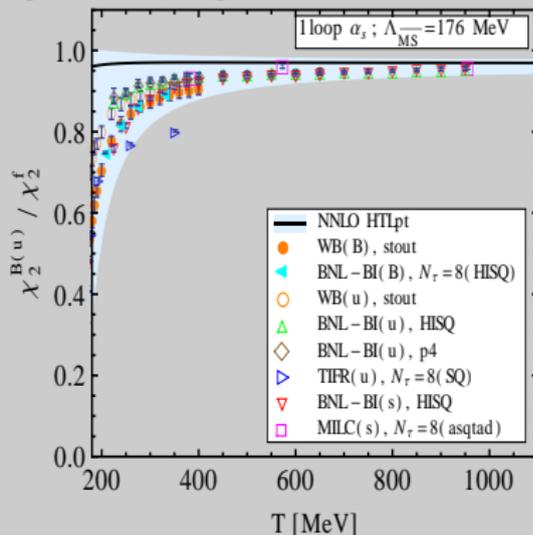
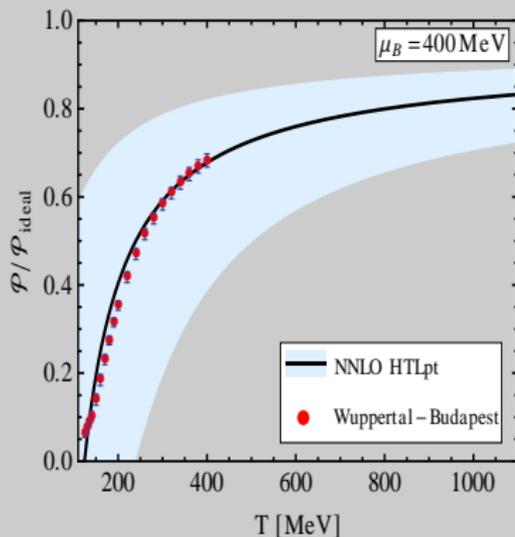
- ☞ Lattice gauge theory has made impressive progress on the calculation of static thermodynamic properties of baryon symmetric QCD matter
- ☞ The equation of state at  $\mu_B = 0$  for physical quark masses is now known with a precision that far exceeds that required in (viscous) hydrodynamics calculations
- ☞ By Taylor expansion pressure for  $\mu_B \neq 0$  is also computed
- ☞ The quasi-critical temperature where chiral suscept. related to chiral symmetry peak has been determined to lie at  $T_c = 155\text{MeV}$

## Scope

- ☞ CP requires further precision in calculation
- ☞ Freeze out has been estimated with a band, requires improvement
- ☞ Dilepton rate and transport coefficients have been computed!

# Thermodynamics and susceptibility in HTLpt

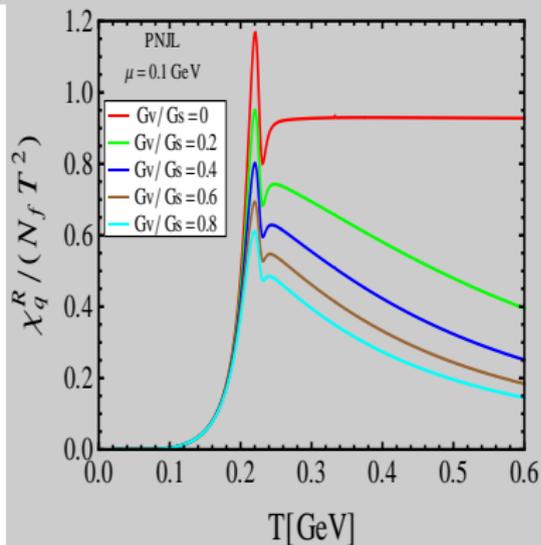
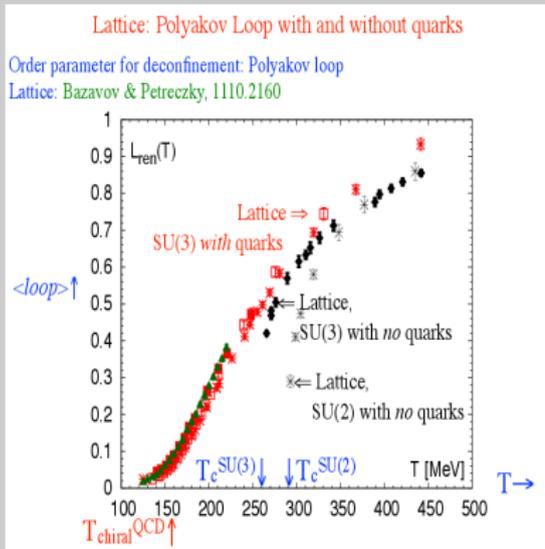
- 3-loop Pressure and susceptibility [N. Haque];



- Trace anomaly, speed of sound and higher order QNS have been computed and agree well with LQCD  $T \geq 250$  MeV
- Scope:** Pressure at large  $\mu_B$  and very small  $T$  for FAIR and Stars; informative to LQCD

# Thermodynamics in Polyakov-Loop-NJL Model

- Pressure and susceptibility [**Aminul**& **Sarbani**; **Stefan Schramm**]



- With(out) Vector interaction ( $G_V$ ) with coupling that doesn't depend on  $T$  and  $\mu_B$
- With vector interaction, pressure and susceptibilities deviate from

## Some more works based on Polyakov Loop

- Finite volume study of thermodynamics [Subrata Sur]
- Evolution of Polyakov loop fields in heavy-ion collisions [P S Saumia]
- Isospin symmetry breaking and baryon-isospin correlation [Sarbani Majumder]
- Phase diagram and CP [Abhijit Bhattacharya; Rajarshi Ray, Sanjay Ghosh]

# Scope

- Temperature dependent vector coupling  $G_V$  is required to match with LQCD data
- $D = 4$  Gluon Condensates in quark and gluon props. [Purnendu Chakraborty: JHEP 1303 (2013) 120]
- Gribov Term in props. [Kharzeev and Levin, arXiv:1501.04622; Su and Tywoniuk, arXiv:1409.3203]

Last two items in above: probably marring the perturbative and nonperturbative approach in  $T_c \leq T \leq 2T_c$

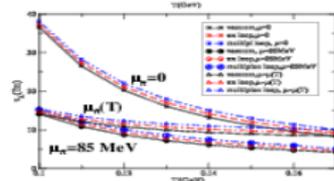
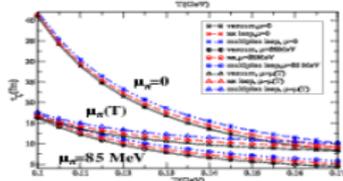
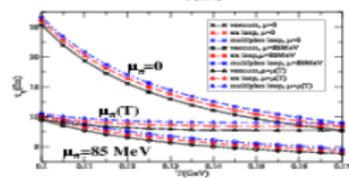
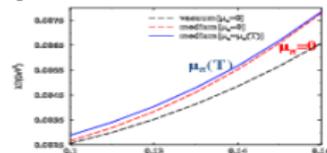
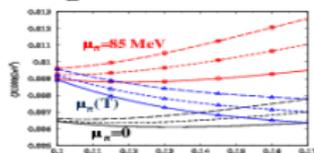
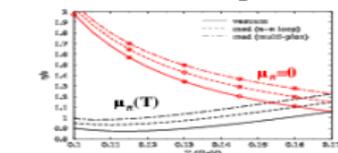
# Transport Coefficients

- They are extracted through vHydro by fitting flow coeffs.
- Also calculated in two ways: Boltzmann Equation & Kubo Formula
- Fitting to data doesn't give  $T$  and  $\mu$  dependence
- Since these are properties of a medium, should  $T$  and  $\mu$  dependence
- So one should use Boltzmann Eq & Kubo formula

# Transport Coefficients

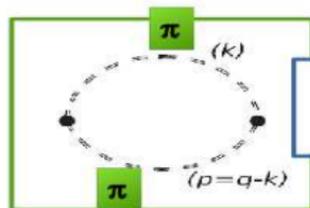
## Bulk and Shear [Sukanya Mitra]

- The **transport coefficients and relaxation time of dissipative flows** have been evaluated from microscopic transport theory for a hot pion gas.
- The effects of a **thermal medium** have been introduced through the **dynamical interaction cross section** that has been estimated from **effective theory at finite temperature**.
- The medium modified interaction cross section gets **suppressed** which in turn **enhances the temperature dependence of the first and second order transport coefficients**.
- A **temperature dependent pion chemical potential** has been used indicating the early chemical freeze out of pion gas which also affects the temperature dependence of transport coefficients.
- Being important inputs of the hydrodynamic evolution equations the **medium modified transport coefficients** are expected to affect the signals extracted from hydrodynamic simulations.



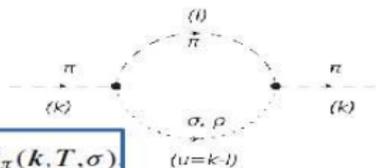
# Transport Coefficients

- $\eta$  [Sabyasachi Ghosh]

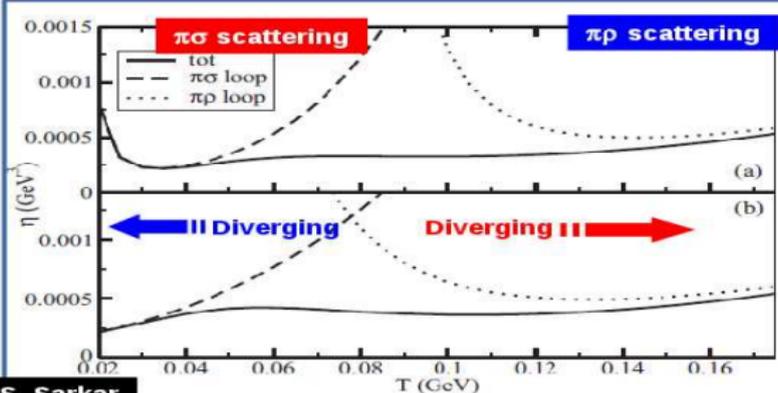


$$\eta_\pi = \frac{\beta}{10\pi^2} \int_0^\infty dk k^6 \frac{d}{\omega_k^2 \Gamma_\pi} n_k (1 + n_k)$$

$$\Gamma_\pi(k, T) = \Gamma_\pi(k, T, \rho) + \Gamma_\pi(k, T, \sigma)$$



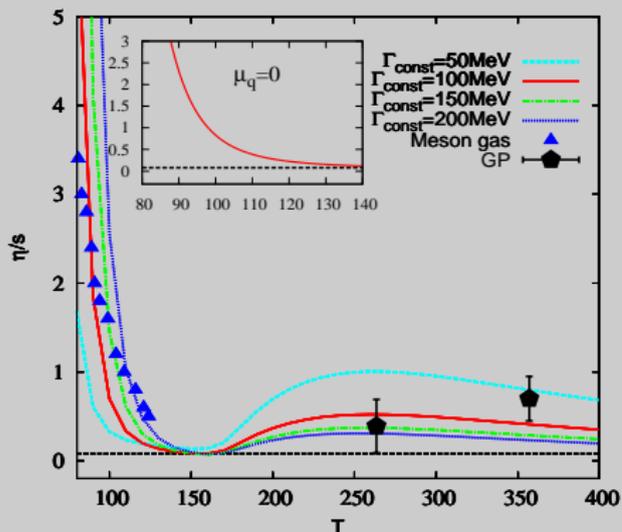
Interestingly, the  $\pi\rho$  and  $\pi\sigma$  contributions play a complementary role in  $\eta$  to be **nondivergent** in the **higher** ( $T > 0.100$  GeV) and **lower** ( $T < 0.100$  GeV) temperature regions, respectively.



S. Ghosh, G. Krein, and S. Sarkar,  
Phys. Rev. C **89**, 045201 (2014).

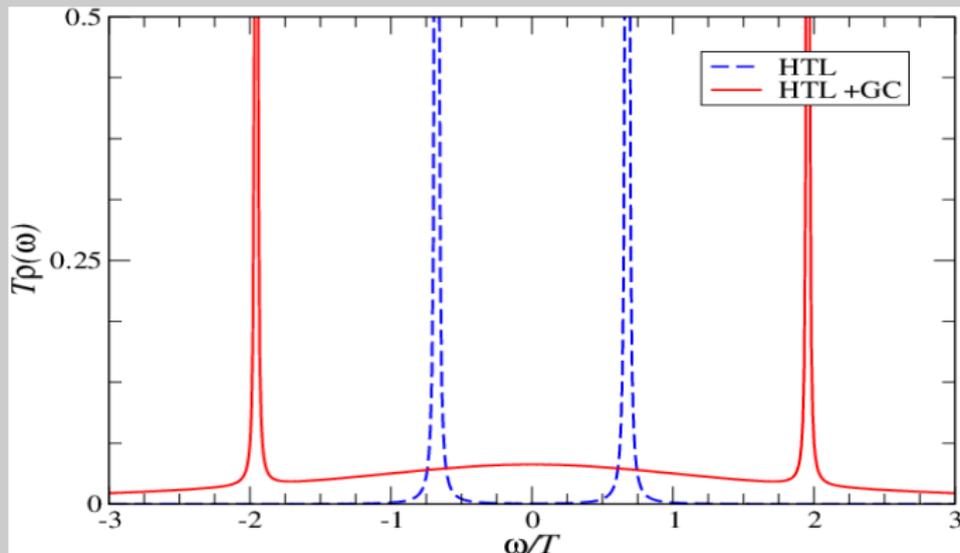
# Transport Coefficients

- $\eta/s$  in PNJL but considering some width in quark propagator [**Sudipa and Kinkar**]



## Scope in Transport Coefficients

- $D = 4$  Gluondensates in Quark Propagator [Purnendu Chakraborty: JHEP 1303 (2013) 120]



- May be with Gribov term too, but not sure at the moment!

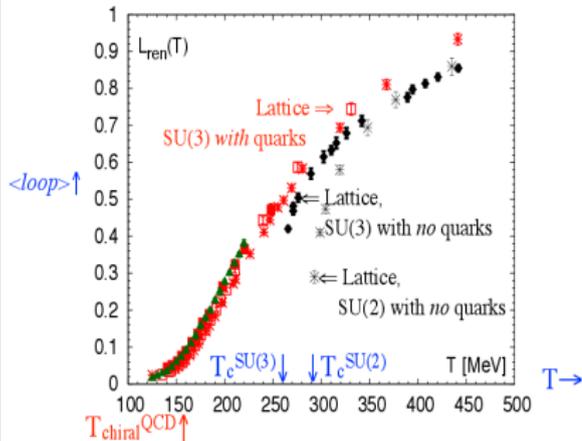
# Dileptons

- Rob Pisarski

Lattice: Polyakov Loop with and without quarks

Order parameter for deconfinement: Polyakov loop

Lattice: Bazavov & Petreczky, 1110.2160



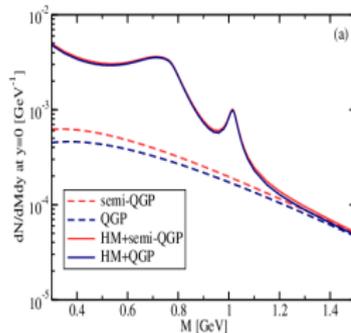
Hydrodynamics: # dileptons

MUSIC: 3+1 hydro @ RHIC:  $\sqrt{s} = 200$  GeV/A, central collisions

Preliminary analysis: only ideal hydro.

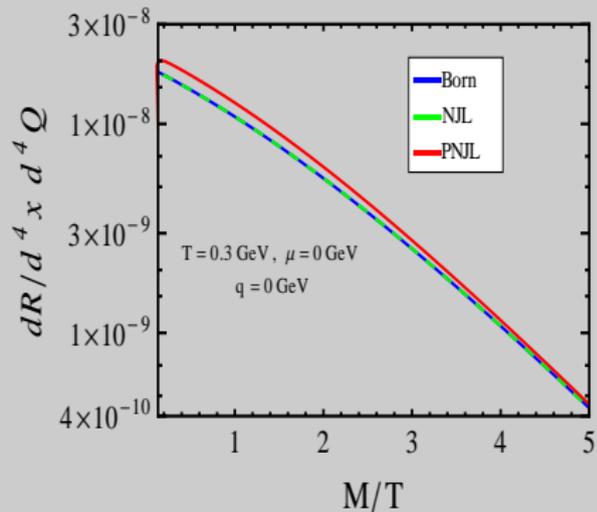
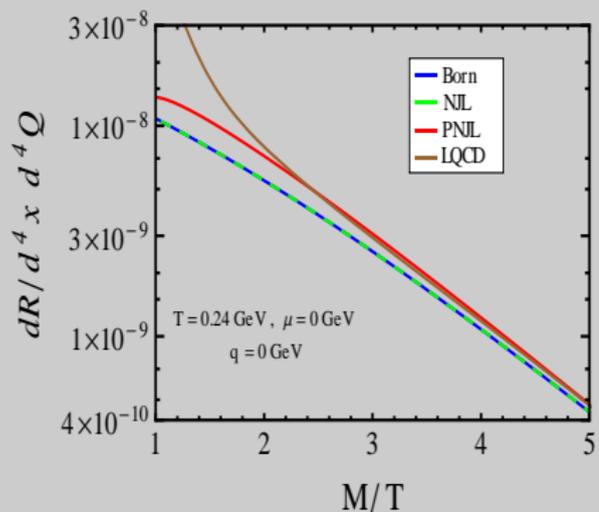
Small enhancement of dileptons in semi-QGP, swamped by hadronic phase.

No matching of semi-QGP to hadronic phase: clearly essential.



# Dileptons

- Aminul and Sarbani



- Dilepton rate with gluon condensates [Aritra & Purnendu]

# Message

- Dilepton Rate is computed in LQCD but bias dependent spectral function
- pQCD (HTL) rate is available for more than a decade
- Nonperturbative rates (with Polyakov Loop, Gluon Condensates) have been computed
- ☞ No match to the enhancement of low mass dilepton spectra
- ☞ At low mass Dileptons from Hadrons essential [Rolf Rapp & H van Hees; Sourav Sarkar]

# Energy Loss and Jet quenching [Abhijit Majumdar, Ivan Vitev]

- Significant Assumptions in E-loss Model:

- The virtuality and splitting of hard parton and how to reduce its virtuality
- The nature of the medium through which the energy parton propagates
- Kinematic approximations for the interaction between medium partons and projectile partons: Eikonal approximation, soft gluon emission, small angle/colinear emission

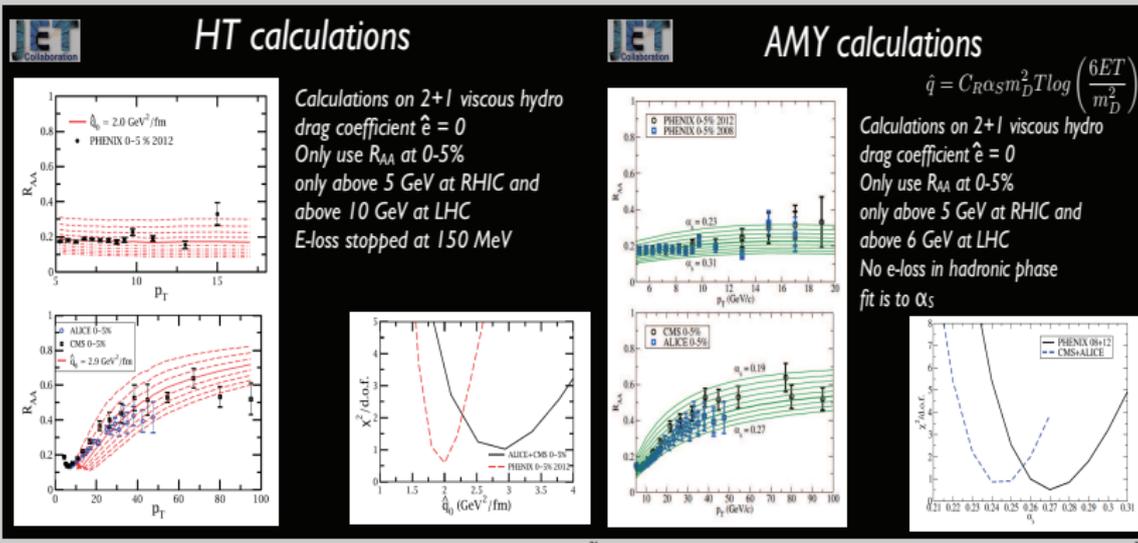
Some extensively used models: BDMPS-Z, GLV, AMY, HT

- Modeling the medium:

- Collection of static scattering centers with specified density distributions
- Thermally equilibrated perturbative medium
- Multigluon exchange between projectile and spatially extended medium in terms of expectation values of higher field operators  $\Rightarrow$  Higher Twist

# Energy Loss and Jet quenching

- Abhijit Mazumdar



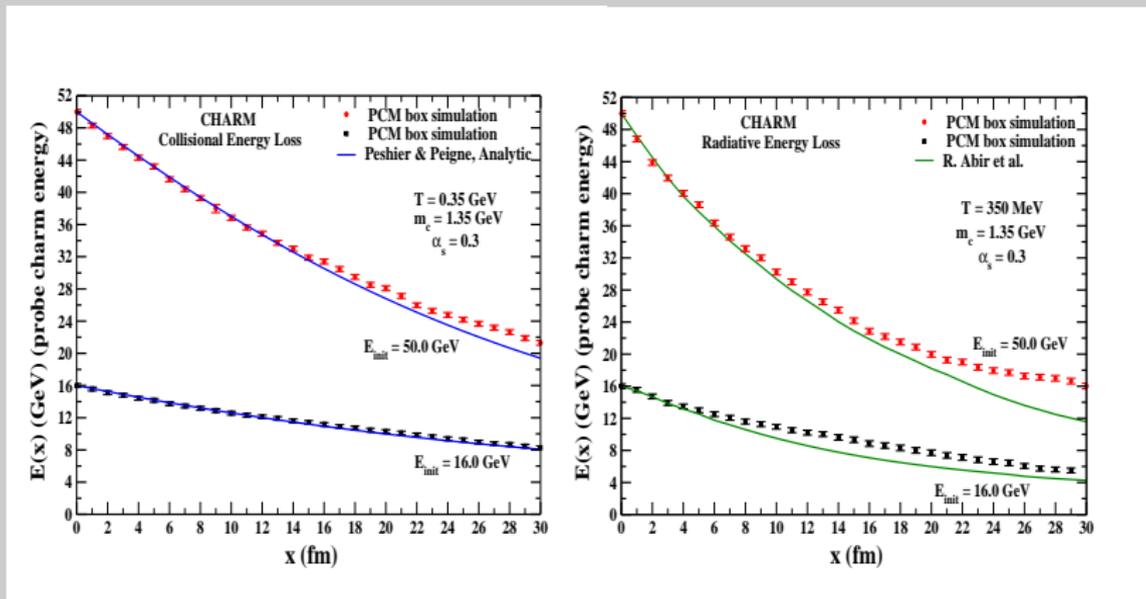
- No collisional  $\hat{e}$  in above; discussed uncertainties in IC of hydro and also in modeling
- Emphasized Importance on collision  $\hat{e}$ ; can be taken into account consistently in a given model!

# Heavy Quark Quenching

- LHC
  - ☞ Charm: LHC  $\sim 10 \times$  RHIC
  - ☞ Bottom: LHC  $\sim 100 \times$  RHIC
- Produced at early time ( $\tau_Q < \tau_{th}$ ); No production at later time
- Produced in pairs (QQ) at early times  $\Rightarrow$  pQCD
- Total no. of HQ gets frozen very early in the history of collisions
- Immediately upon their production they will propagate through QGP
- One is left with the task of determining the HQ distribution
- Details of the distribution may reflect the characteristics and development of QGP
- Generalized Dead-Cone formula for radiative E-loss
- Fokker-Planck Eq. distribution for collisional & radiative loss
- $R_{AA}$ : D-meson spectra; semileptonic decay of Charmed and Bottomed Meson
- HQ and LQ lose energy in similar fashion

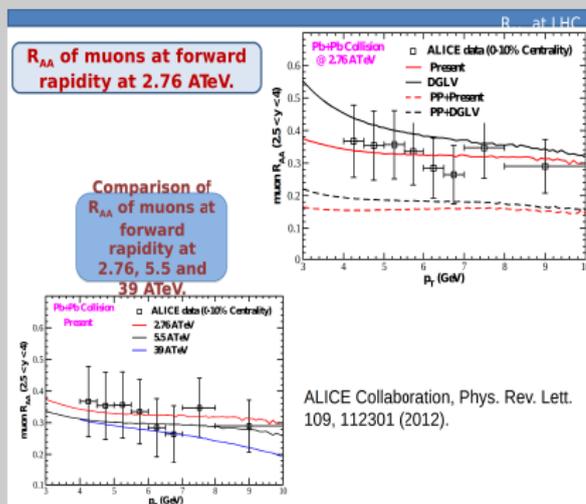
# Heavy Quark Quenching

- Charm Evolution: Analytic model vs PCM(VNI/BMS) [Md. Younus]



# Heavy Quark Quenching

- U. Jamil



- K. Saraswat; Souvik P Addya

# Quarkonia Melting

- Color Screening



## Quarkonia and the concept of suppression

**Charmonia:**  $J/\psi$ ,  $\Psi'$ ,  $\chi_c$       **Bottomonia:**  $\Upsilon$  (1S),  $\Upsilon$  (2S),  $\Upsilon$  (3S)  
 Heavy quarks carry the information of early stage of collisions:

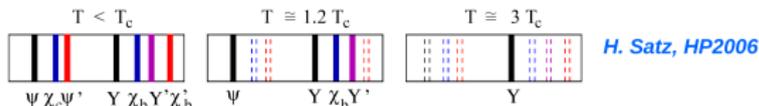
- Charm and bottom quarks are massive.
- Formation takes place early only in the collision.

### Proposed Signature: QCD phase transition

Color screening of static potential between heavy quarks:

- $J/\psi$  suppression: Matsui and Satz, *Phys. Lett. B* **178** (1986) 416
- Suppression of states is determined by  $T_c$  and their binding energy
- Lattice QCD: Evaluation of spectral functions  $\Rightarrow T_{\text{melting}}$

**Deconfinement  $\rightarrow$  Color screening  $\rightarrow$  heavy quarkonia states “dissolved”**



$$T_{\text{diss}}(\Psi') \approx T_{\text{diss}}(\chi_c) < T_{\text{diss}}(\Upsilon(3S)) < T_{\text{diss}}(J/\psi) \approx T_{\text{diss}}(\Upsilon(2S)) < T_{\text{diss}}(\Upsilon(1S))$$

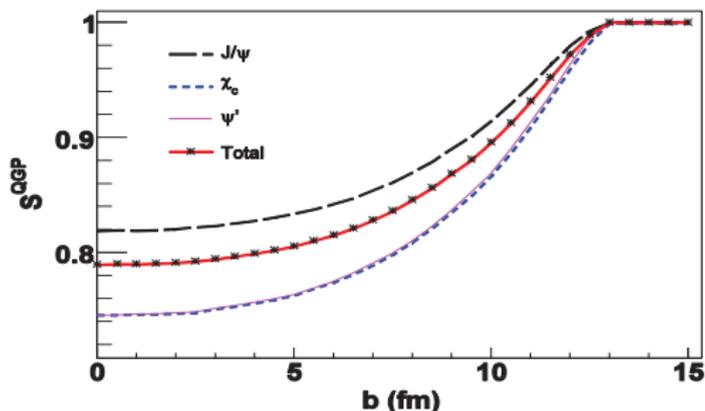
Their suppression pattern is a **thermometer** of the QCD matter

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# Quarkonia Melting

- P. Bhaduri
- Color screening
- Described pA with cold nuclear effects
- Prediction for FAIR

## Centrality dependence of inclusive survival probability

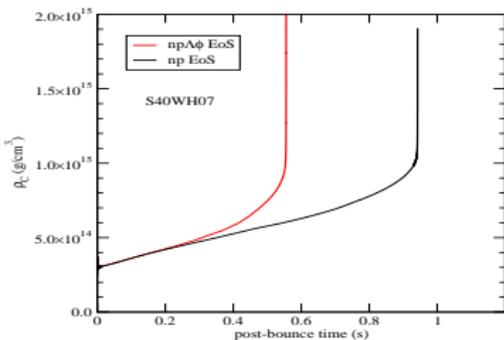
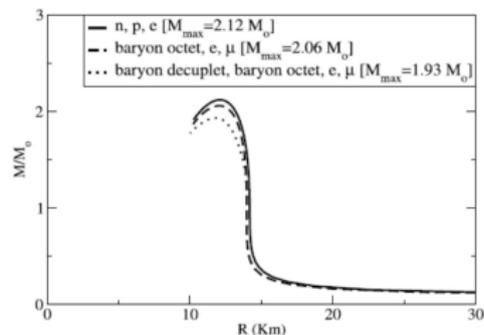


Effects of feed down is included: 30 % from  $\chi$  and 10 % from  $\psi'$

$$S_{\text{tot}} = 0.6S_{J/\psi} + 0.3S_{\chi} + 0.1S_{\psi'}$$

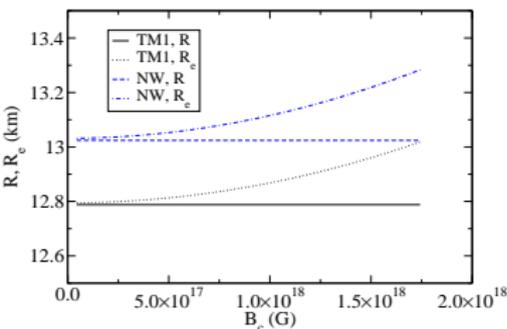
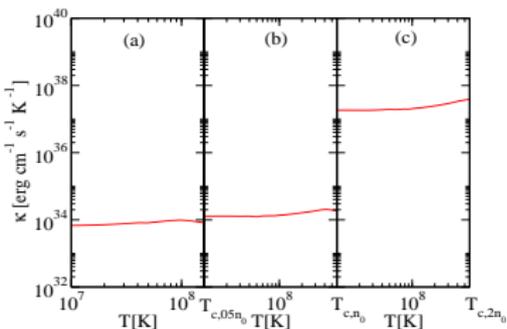
# Does exotic matter exist in $2 M_{\odot}$ Neutron Stars?

Accurately measured highest Neutron Star of  $2.01 \pm 0.04 M_{\odot}$  puts stringent conditions on the EoS. [ J. Antoniadis et al., Science 340 (2013)]



- ▶ **Hyperon puzzle - Softer EoS, smaller maximum mass** (S. Schramm's talk)
- ▶ **Chiral EoS including hyperons,  $\Delta$  resonances** (V. Dexheimer, S. Schramm, ApJ683 (2008))
- ▶ **Vector meson repulsion makes the EoS stiffer resulting in higher maximum mass**
- ▶ **Hadron look-alike quark matter EoS leads to heavy neutron stars.** (M. Alford et al. ApJ 629 (2005))
- ▶ **First  $\Lambda$  hyperon SN EoS compatible with  $2 M_{\odot}$  NS** (S. Banik, M. Hempel, D. Bandyopadhyay, ApJS 214 (2014))
- ▶ **The supernova simulation with this hyperon EoS leads to early collapse of the protoneutron star to a black hole** (P. Char's talk)

# Superfluidity and Strong Magnetic Fields in NS



- ▶ Rapid cooling of the Neutron Star in Cas A is attributed to the neutron superfluidity in the core [ D. Page et al., PRL106 (2011)]
- ▶ The contribution of superfluid phonons to the thermal conductivity computed solving Boltzmann equation are shown here (S. Sarkar's talk)
- ▶ Neutron Stars are laboratories of dense matter under extreme physical conditions. The surface magnetic field could be as high as  $10^{16}$  G.
- ▶ Such high magnetic fields influence the EoS through Landau quantization
- ▶ Anisotropic magnetic pressure deforms the neutron star (R. Mallick's talk)

# Quark Nuggets with $Z(3)$ Symmetry

## ▶ Abhishek Atreya

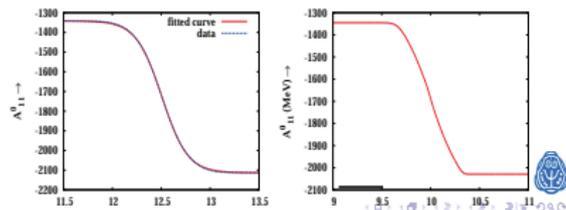
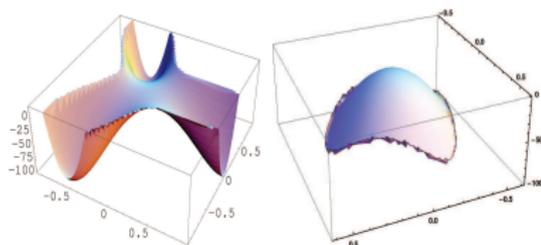
quarks may be trapped in  $Z(3)$  domain

Quark Nuggets From  $Z(3)$  Walls

CP Violation and Anti-Nuggets

Quark Nuggets From  $Z(3)$  Walls

CP Violation and Anti-Nuggets



## Baryon Density Profile

$$\rho(R) = \frac{\dot{N}_i}{4\pi v_W R^2}$$

- $n_b \sim 10^{52} - 10^{53}$  for  $R < 1$  m.

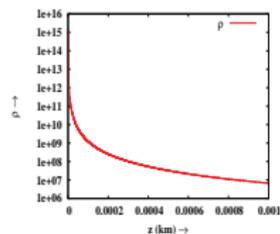


Figure : Baryon density left behind by collapsing wall.