Compton back-scattering photons for QGP tomography

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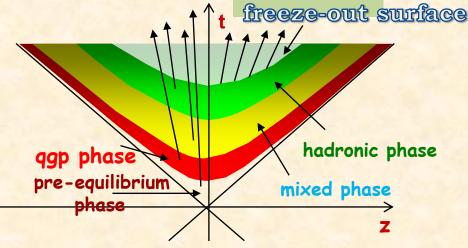
Ref: Phys. Rev. C 90, 034911 (2014)

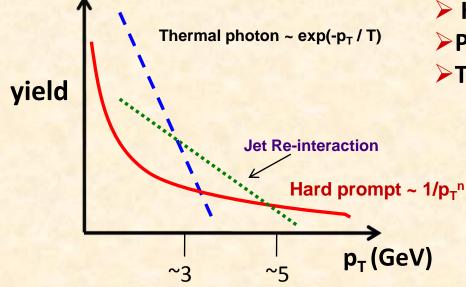
CNT QGP Meet-2015, 16-20 Nov, Kolkata

Electromagnetic probes in Heavy Ion Collisions:

✓ Direct photons are considered as the most cleanest probe in HIC

 \checkmark Due to weak coupling ($\alpha_{\rm e}/\alpha_{\rm s}\sim$ 10⁻²) with medium, direct photons carry information undistorted from each stages of evolution



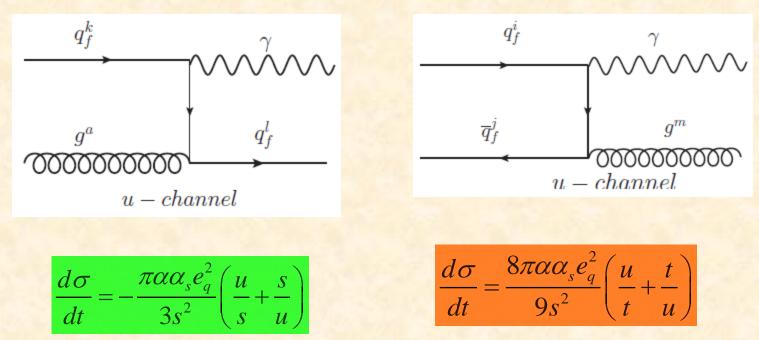


Hard photons (Direct + Jet fragment)
 Pre-equilibrium and jet-medium photons
 Thermal photons (QGP + Hadron matter)

Experimental challenge to separate different sources of direct photon

Photons from re-scattering of jets in quark gluon plasma

QCD Compton and Annihilation process :



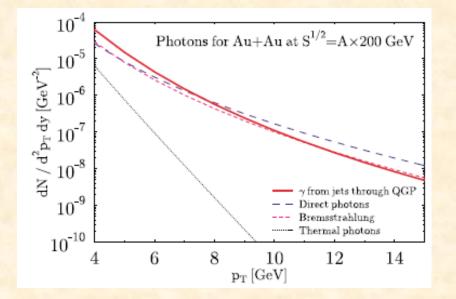
Cross sections are maximum for small values of t and u Backward scatt. $\vec{p}_{\gamma} \approx \vec{p}_q(\vec{p}_{jet})$ $\vec{p}_{\gamma} \approx \vec{p}_{\bar{q}}(\vec{p}_{jet})$

Compton back-scattering is often used for the production of high energy laser beam

Jet-photon contd..

Total inclusive yield:

$$E_{\gamma} \frac{dN}{d^4 x d^3 p_{\gamma}} = \frac{\alpha \alpha_s}{4\pi^2} \sum_{f} (\frac{e_f}{e})^2 [f_q(x, p_{\gamma}) + f_{\bar{q}}(x, p_{\gamma})] T^2 [\ln \frac{3E_{\gamma}}{\alpha_s \pi T} + C]$$



First proposed by : Fries, Muller, Srivastava (PRL,90, 132301 (2003))

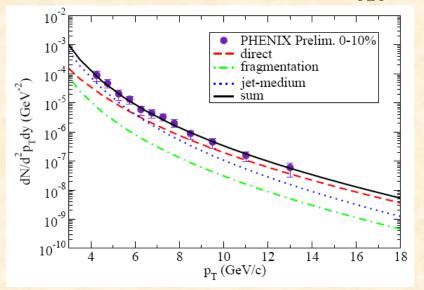
Shows substantial contribution for $p_T \le 6$ GeV at RHIC energies

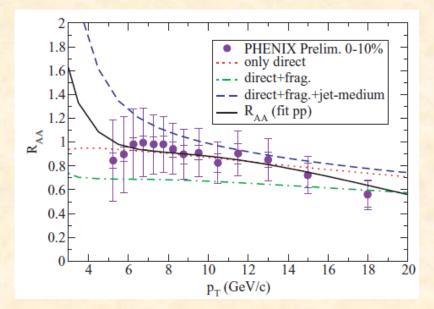
Experimental measurement of photons :

Inclusive yield and Nuclear modification factor of direct photons

Azimuthal momentum anisotropy coeff. (v₂)

Inclusive yield and R_{AA}:





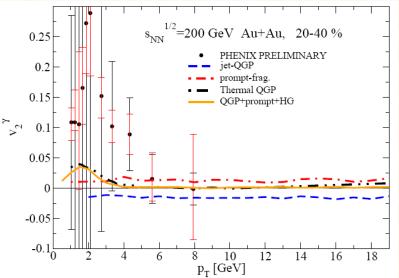
Qin, Ruppert, Gale, Jeon & Moore, PRC 80 (2009)

Azimuthal momentum anisotropy:

- •Jet-medium photons shows negative v₂.
- Theoretical predictions are inconclusive.

Chatterjee, Frodermann, Heinz, Srivastava; PRL 96 (2006) Turbide, Gale, Fries ; PRL 96 (2006)

Not promising, so far...



Turbide et al. PRC 77 (2008)

Jet - tagged photon measurement:

Motivation:

*The back-scattered photons has strong correlation with the parent jet momentum

✤ Jets are produced back-to back in the medium

Strategy:

Fix the momentum and rapidity of the away-side leading jet

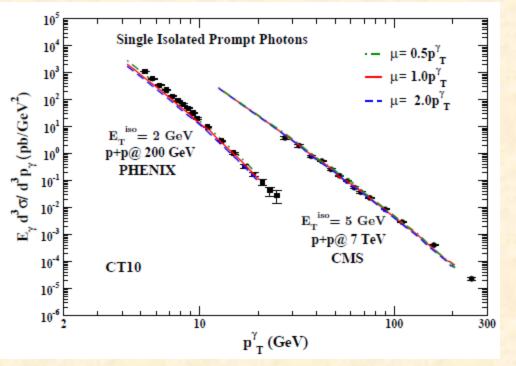
> Study the photons, very close to the away-side jet axis

>The initial hard photons are treated as background

>Get rid of thermal and pre-equilibrium photons



Estimation of Background



SD, Pramana 82 (2014)

JETPHOX: S.Catani, M. Fontannaz, J. Ph. Guillet, E. Pilon JHEP 05 (2002) 028

Photons from initial hard collision + Fragmentation of jets

The background is calculated from the NLO package JETPHOX

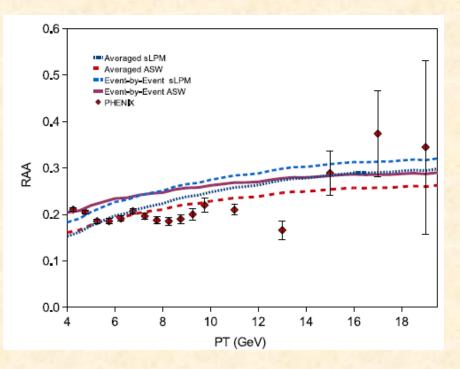
EPS09 nuclear pdf is used for A+A collisions

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•No E<sub>T</sub> cut for A+A case
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The energy loss model : ppm

We have used a longitudinally expanding, boost-invariant fireball model by: Rodriguez, Fries, Ramirez PLB 693 (2010)

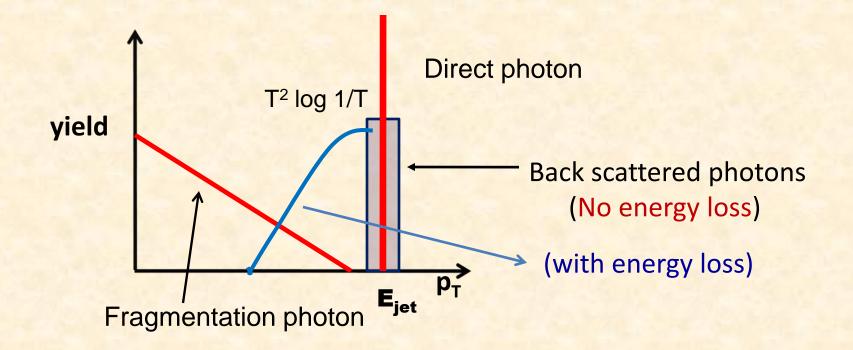
Path travelled by the jet: $I_{\beta}(\vec{r}, \varphi) = \int d\tau \tau^{\beta} \rho(\vec{r} + \tau \vec{e}_{\varphi})$



We have used LPM type of energy loss; $\beta = 1$ $\Delta E = C_{LPM} I_1(r, \varphi)$

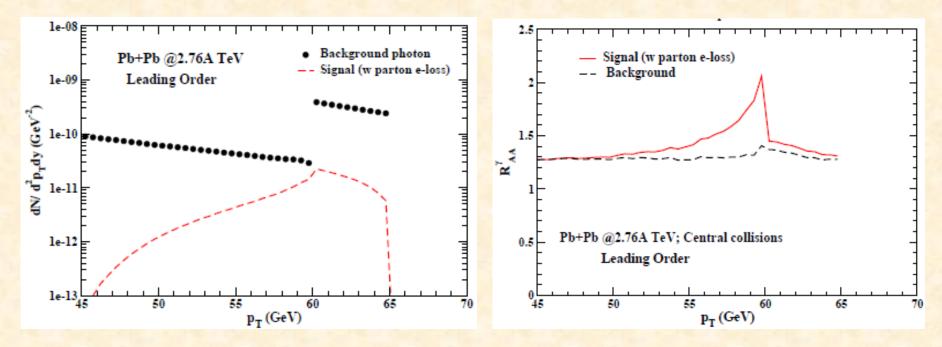
The coeff. C_{LPM} is determined from the fitting of R_{AA} of hadrons

Jet-tagged photons at Leading order: a schematic view



Important Information: Temperature of the medium Energy loss of partons before back-scattering QCD back-scattering photons at LHC : LO R_{AA} : (Signal + Background)_{AA} / N_{coll} ×(Background)_{pp} •For central Pb+Pb collisions at 2.76 TeV at mid rapidity

- Photons opposite to the 60-65 GeV jet within ± 15 degrees
- The quarks suffer energy loss before conversion



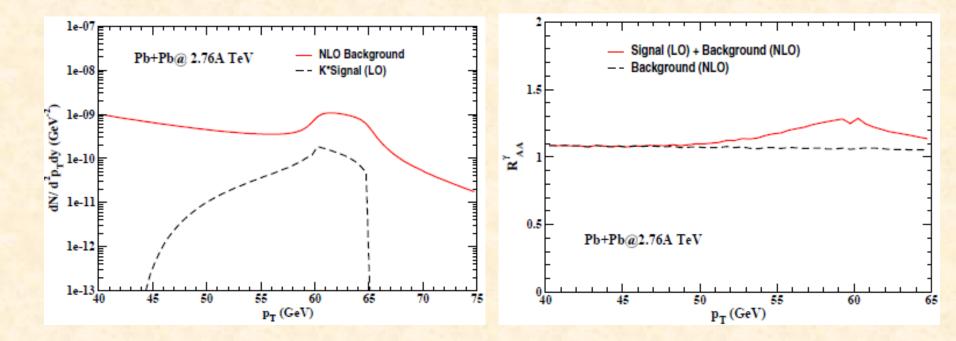
A clear back scattering peak can be seen just below trigger window

QCD back-scattering photons at LHC : NLO

Back ground is calculated in the Next-to Leading order

Kinematics of jet-conversion is still leading order

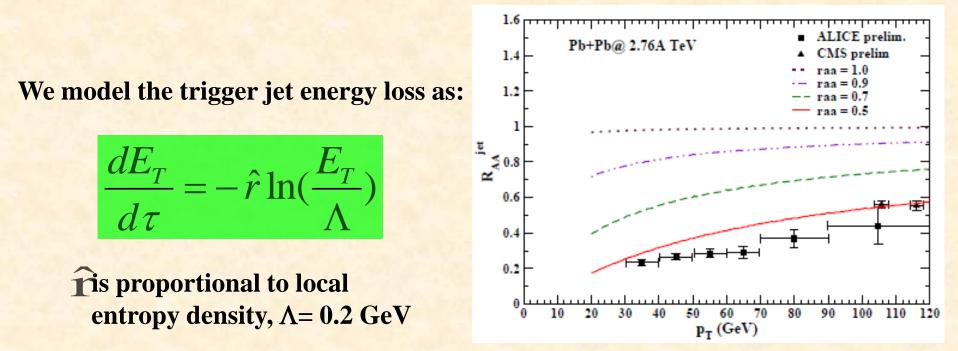
Parton energy loss is accounted





Effect of trigger jet energy loss

Recent measurements at LHC (arXiv:1304.5945, PLB 2015) suggest a strong suppression of trigger jets in central collisions.



SD, R J Fries, D K Srivastava PRC 90 (2014)

Trigger jet energy loss affects the Background (direct and fragmentation photon) and as well as the Signal

The parton-jet pair distribution:

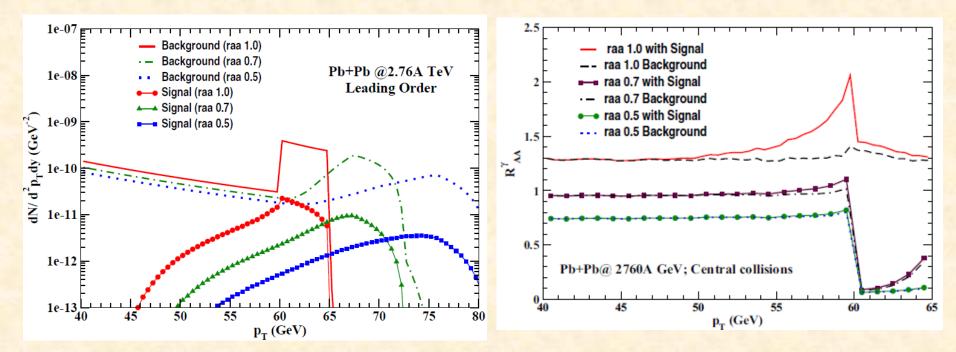
The parton-jet pair distribution integrated over a window Γ in $E_T - \phi - \eta$ space

$$\begin{split} f_q^{\mathcal{T}_j}(\mathbf{p}_q, x) &= \frac{(2\pi)^3}{g_q \tau p_T} \delta(y - \eta) \rho(\tau, \mathbf{x}_{\perp}^0) \\ \times \int_{\mathcal{T}_j} dE_T \, dy_j d\phi_j E_q \frac{dN}{d^3 p_q dE_T dy_j d\phi_j} \Big|_{\substack{\mathbf{p}_q^0 = \mathbf{p}_q + \Delta \mathbf{p}_q \\ E_T^0 = E_T + \Delta E_T}} \end{split}$$

The parton-jet pair are evolved through the medium while their respective energy losses are also accounted.

> The parton back-scattering probability is also computed along the way.

Background and Signal: w trigger jet energy loss



SD, R J Fries, D K Srivastava PRC 90 (2014)

Signal and Background are calculated for Leading order kinematics

Both trigger jet and parton energy loss are taken into account

Trigger jet energy loss tends to wash out the strong correlation with parent jet

Summary & Conclusion:

- •Jet-medium back scattering photon is an important signature of thermalized matter created in relativistic heavy ion collisions.
- •We propose the use of trigger jet to identify this particular source of direct photons.
- •Jet-medium photons shows characteristic enhancement in the nuclear modification factor of direct photon production at large momentum.
- •The peak is clearly visible in Leading order but weakens for the radiative corrections to the process and trigger jet energy loss.
- •The shift of the peak from the trigger jet window provides complimentary measure of parton energy loss in the medium.
- •Separation of this signal from other photon sources depends crucially on the initial trigger jet energy estimation.



Jet-Photon Conversion

The rate of production:

$$E_{\gamma} \frac{dN^{(A)}}{d^4 x d^3 p_{\gamma}} = \frac{16E_{\gamma}}{2(2\pi)^6} \sum_{q=1}^{N_f} f_q(p_{\gamma})$$

$$\times \int d^3 p f_{\overline{q}}(p) [1 + f_g(p)] \sigma^{(A)}(s) \frac{\sqrt{s(s - 4m^2)}}{2E_{\gamma}E} + (q \leftrightarrow \overline{q})$$

$$E_{\gamma} \frac{dN^{(C)}}{d^4 x d^3 p_{\gamma}} = \frac{16E_{\gamma}}{2(2\pi)^6} \sum_{q=1}^{N_f} f_q(p_{\gamma})$$

$$\times \int d^3 p f_g(p) [1 - f_q(p)] \sigma^{(C)}(s) \frac{s - m^2}{2E_{\gamma}E} + (q \leftrightarrow \overline{q})$$

$$f_g^q(p) = f_g^{jet}(p) + f_g^{th}(p)$$

$$dN^{\gamma} = f_{th} \otimes f_{th} + f_{th} \otimes f_{jet} + \dots \text{ Jet-converted photons}$$

Results at RHIC energy

